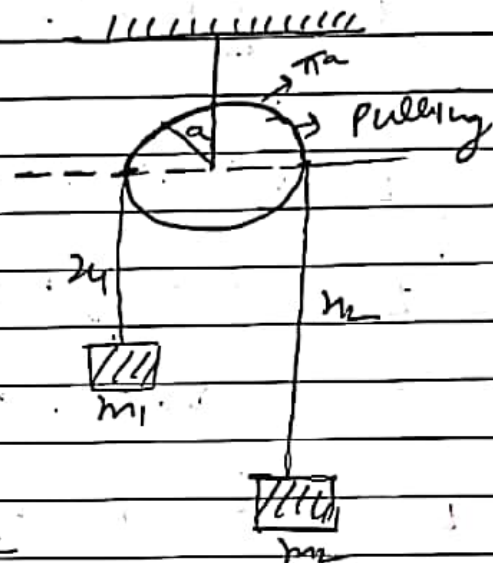


m-imp

\* At Woods machine of

Let us consider two masses  $m_1$  &  $m_2$  which are hanging on the wire by a pulley. And pulley is frictionless.



And there are two coordinates  $x_1$  &  $x_2$  the motion of these particles by constraint condition can be written as.

$$L = x_1 + x_2 + \pi a$$

$$x_2 = L - x_1 - \pi a \quad \text{--- (1)}$$

diff. w.r to  $t$

$$\dot{x}_2 = -\dot{x}_1 \quad \text{--- (2)}$$

Where  $L$  is the length of the wire and  $a$  is the Radius of pulley.

Now the K.E of both masses can be written as.

$$T = \frac{1}{2} m v^2 = \frac{1}{2} m_1 \dot{x}_1^2 + \frac{1}{2} m_2 \dot{x}_2^2 \quad \text{--- (3)}$$

$$V = V_1 + V_2$$

due to the gravitational force the P.G of both masses should be -ve.

$$Mz - m_1 g x_1 = m_2 g x_2 \quad \text{--- (1)}$$

we know that

$$L = T - V$$

$$L = \frac{1}{2} m_1 \dot{x}_1^2 + \frac{1}{2} m_2 \dot{x}_2^2 + m_1 g x_1 + m_2 g x_2$$

from eq (1)

$$L = \frac{1}{2} m_1 \dot{x}_1^2 + \frac{1}{2} m_2 \dot{x}_1^2 + m_1 g x_1 + m_2 g (L - x_1 - \pi a)$$

$$L = \frac{1}{2} \dot{x}_1^2 [m_1 + m_2] + g x_1 (m_1 - m_2) + m_2 g (L - \pi a)$$

$$V_0 = -m_2 g (L - \pi a)$$

$$L = \frac{1}{2} \dot{x}_1^2 [m_1 + m_2] + g x_1 (m_1 - m_2) + V_0$$

now from the Lagrange Hamilton eq:

$$\frac{\partial L}{\partial x_1} - \frac{d}{dt} \left( \frac{\partial L}{\partial \dot{x}_1} \right) = 0$$

$$\frac{d}{dt} \left( \frac{\partial L}{\partial \dot{x}_1} \right) - \frac{\partial L}{\partial x_1} = 0$$

diff w.r to  $x_1$  and  $\dot{x}_1$  we get

$$\frac{\partial L}{\partial x_1} = g (m_1 - m_2)$$

$$\frac{\partial L}{\partial \dot{x}_1} = \dot{x}_1 (m_1 + m_2)$$

$$\frac{d}{dt} \left[ \dot{x}_1 (m_1 + m_2) - g (m_1 - m_2) \right] = 0$$

$$\therefore \dot{x}_1 (m_1 + m_2) - g (m_1 - m_2) = 0$$

$$\dot{x}_1 (m_1 + m_2) = g (m_1 - m_2)$$

$$\dot{x} = \frac{g (m_1 - m_2)}{m_1 + m_2}$$

$$m_1 + m_2$$

Imp  
\* D'Alembert's principle:

This principle is based upon the principle of virtual work.

$$\text{i.e. } W = \sum_{i=1}^N \vec{F}_i \cdot \delta \vec{x}_i = 0 \quad \text{--- (1)}$$

If there is small  $N$  particles then each particle exerts the force  $F_1, F_2, \dots, F_N$  and gets the displacements  $\delta \vec{x}_1, \delta \vec{x}_2, \dots, \delta \vec{x}_N$ .  
If the particle is restricted i.e.

the constraints condition is applied on a particle then the total force of the body is equal to the sum of applied force

$$\vec{F}_i = \vec{F}_{ia} + f_i \quad \text{--- (2)}$$

[  $F_{ia}$  = applied force  
 $f_i$  = constraints force ]

Now we know that the conservative force is always applied when the displacement is restricted i.e. the constraints force is  $\perp$  to the displacement

$$\sum_{i=1}^N \vec{F}_{ia} \cdot \delta \vec{x}_i = 0 \quad \text{--- (3)}$$

The equation is termed as virtual work.

To interpret to equilibrium system

D'Alembert adopted an idea of the reversed force

He conceived that a system will remain in equilibrium under the action of a force equal to the



Actual force  $\vec{F}_i$  And Reverse Effective force  $\vec{P}_i$

$$\vec{F}_i = \frac{d\vec{P}_i}{dt} = \vec{P}_i$$

$$\vec{F}_i = \vec{P}_i$$

$$\vec{F}_i - \vec{P}_i = 0$$

$$\sum_C^N (\vec{F}_i - \vec{P}_i) \cdot \delta \vec{x}_i = 0$$

$$\sum_C^N (\vec{F}_i - \vec{P}_i) \cdot \delta \vec{x}_i + \sum_C^N \vec{f}_i \cdot \delta \vec{x}_i = 0$$

Since force of constraints are no more in picture, so it is better to drop.

$$\sum_C^N (\vec{F}_i - \vec{P}_i) \cdot \delta \vec{x}_i = 0$$

Which is called D'Alembert's Principle which shows that the Reverse Effective force always decrease the value of Applied force